

### First-Order Transition in Doped Polyacetylene

Recently Kivelson and Heeger (KH)<sup>1</sup> have proposed that doped polyacetylene undergoes a first-order transition from a lattice of charged solitons into a polaron metal. In this Comment I show that thermodynamic consequences of this conjecture explain and tie together a variety of experimental data. The experiments considered are (a) a very sharp transition in the ESR response as function of the electrochemical potential  $\mu$ ,<sup>2</sup> (b) a plateau in the relation  $\mu(\rho)$ ,<sup>3</sup> where  $\rho$  is the excess charge density on the polyacetylene chains (or the doping density), and (c) the diversity of data on the Pauli susceptibility  $\chi_P$  for chemically doped samples.<sup>2,4</sup>

The KH conjecture is shown schematically in Fig. 1(a). The free energy is monotonic, but has a cusp at some density  $\rho_c$  where the free energies of two distinct configurations cross. KH propose that  $\rho < \rho_c$  corresponds to charged solitons, while  $\rho > \rho_c$  corresponds to polarons. The following reasoning, however, is independent of the detailed nature of the two configurations.

The form  $F(\rho)$  in Fig. 1(a) implies that there are two densities  $\rho_1, \rho_2$  connected by a line  $F_{eq}(\rho)$  [the dashed line in Fig. 1(a)] with slope  $\mu_c$ ,

$$F_{eq}(\rho) = F(\rho_1) + \mu_c(\rho - \rho_1), \quad \rho_1 < \rho < \rho_2, \quad (1)$$

such that the line is tangent to  $F(\rho)$  at both  $\rho_1$  and  $\rho_2$ . In terms of the chemical potential  $\mu = \partial F / \partial \rho$  [Fig. 1(b)],  $\mu_c = \mu(\rho_1) = \mu(\rho_2)$ . Equation (1) is equivalent to the finding of a  $\mu_c$  for which the Maxwell equal-area condition holds,

$$\int_{\rho_1}^{\rho_2} \mu(\rho) d\rho = \mu_c(\rho_2 - \rho_1). \quad (2)$$

The meaning of this standard construction is phase separation: Divide the total volume  $V$  into volumes  $V_1, V_2$  with densities  $\rho_1, \rho_2$ , respectively, such that  $V = V_1 + V_2$  and  $\rho V = \rho_1 V_1 + \rho_2 V_2$ . The resulting free energy  $[F(\rho_1)V_1 + F(\rho_2)V_2]/V$  is precisely Eq. (1), which is lower than the homogeneous  $F(\rho)$  [Fig. 1(a)]. Thus for  $\rho_1 < \rho < \rho_2$ , the equilibrium situation is that of a two-phase system separated in space. It also follows that it is  $\mu_c$  (and *not*  $\rho_c$ ) which defines a proper first-order phase transition; at  $\mu_c$ , the density jumps from  $\rho_1$  to  $\rho_2$ .

It is interesting, however, that the KH conjecture implies  $\partial\mu/\partial\rho > 0$ , so that the homogeneous solution is locally stable. When this metastable situation is realized, it should exhibit hysteresis around  $\rho_c$ .

Phase separation in doped polyacetylene is unique since one phase has  $\chi_P = 0$  (solitons), while the other phase has  $\chi_P \neq 0$  (polarons). Thus  $\chi_P(\mu)$  shows an abrupt jump at  $\mu_c$  in agreement with recent data.<sup>2</sup> The

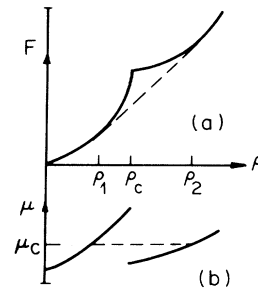


FIG. 1. (a) Schematic free energy for a homogeneous system (full lines) and a phase separated system (dashed line); (b) as in (a) for the chemical potential.

plateau in  $\mu(\rho)$  (Ref. 3) is also an obvious consequence [Fig. 1(b)].

When  $\rho$  is fixed and  $\rho_1 < \rho < \rho_2$ , phase separation predicts that  $\chi_P$  increases with the volume of the  $\rho_2$  phase, i.e.,  $\chi_P(\rho) \sim \rho_2 V_2 / V = \rho_2(\rho - \rho_1) / (\rho_2 - \rho_1)$ . Phase separation also accounts for the presence of two types of spins.<sup>4</sup> It is possible, however, to achieve the homogeneous metastable phase with  $\chi_P = 0$  for  $\rho \leq \rho_c$ . The latter situation is of considerable interest in view of its unusual relation with the conductivity.<sup>1,2</sup> Metastability thus accounts for the diversity of the  $\chi_P(\rho)$  data.<sup>2,4</sup>

Data with Na doping<sup>3</sup> shows  $\rho_1 \approx 1\%$  and  $\rho_2 \approx 6\%$ , which correlates well with the large jump in  $\chi_P(\mu)$ .<sup>2</sup> Similar data on other dopants and a reanalysis of  $\chi_P(\rho)$  data should establish the important values of  $\rho_1, \rho_2$ . The transition is driven by on-chain interactions to the extent that  $\rho_1, \rho_2$  are dopant independent.

Finally, it should be noticed that the microscopic origin of the KH conjecture and the nature of the polaron metal are not conclusive.<sup>1</sup> On the other hand, the KH conjecture does account, as shown here, for a large variety of experimental data.

Baruch Horovitz  
Department of Physics  
Ben-Gurion University  
Beer-Sheva 84105, Israel

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